

Measurement of self-generated spontaneous fields and their effects on ICF ion kinetic dynamics

National Ignition Facility

LaB Conference, July 22-28, 2017
Moscow - St Petersburg, Russia

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Collaborators

**C. K. LI¹, F. H. Séguin¹, J. A. Frenje¹, R. D. Petrasso¹, P.-E. Masson-Labprde²,
S. Laffite², V. Tassin², P. A. Amendt³, H. G. Rinderknecht³, S. C. Wilks³, N. M.
Hoffman⁴, A. B. Zylstra⁴, S. Atzeni⁵, R. Betti⁶, M. Rosenberg⁶ and T. C. Sangster⁶**

¹ *Massachusetts Institute of Technology*

² *CEA*

³ *Lawrence Livermore National Laboratory*

⁴ *Los Alamos National Laboratory*

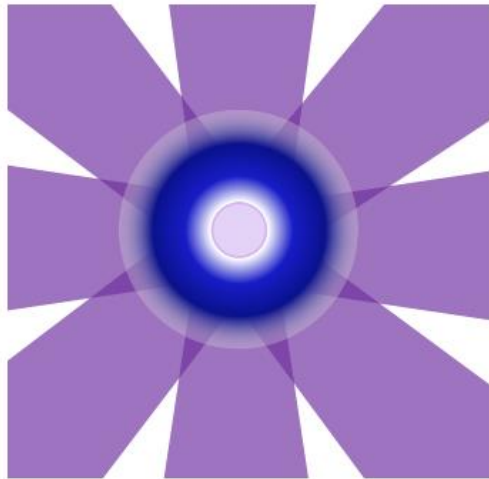
⁵ *University of Rome*

⁶ *University of Rochester*

Two approaches to ICF implosions

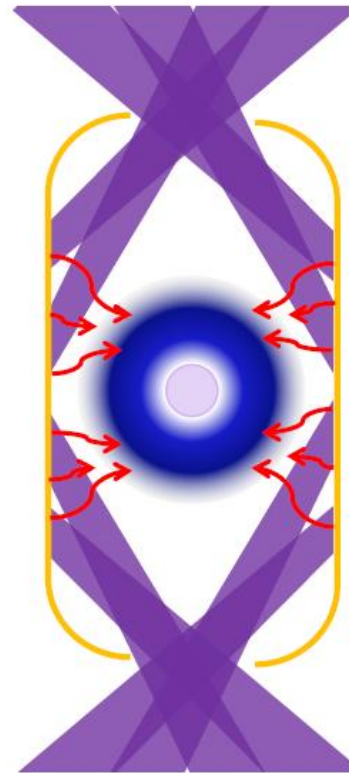
Direct-drive (OMEGA)

Laser directly
irradiates capsule

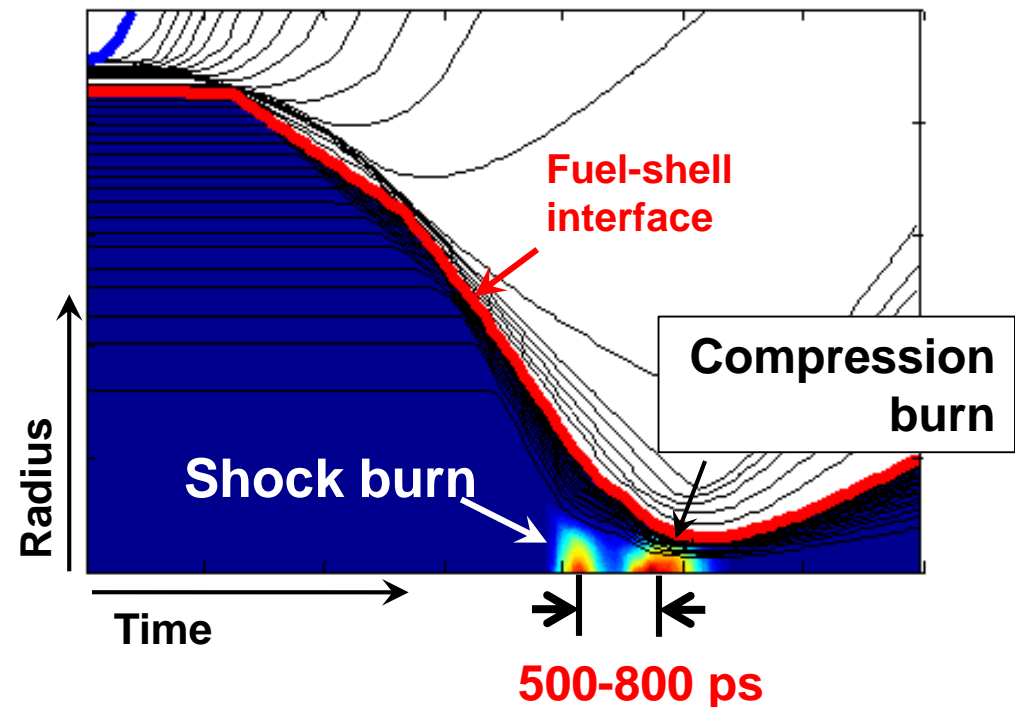
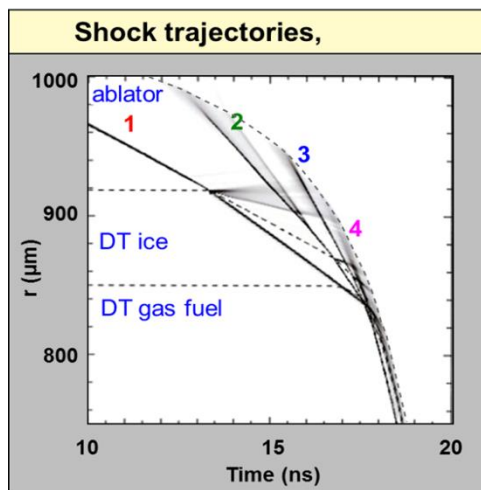
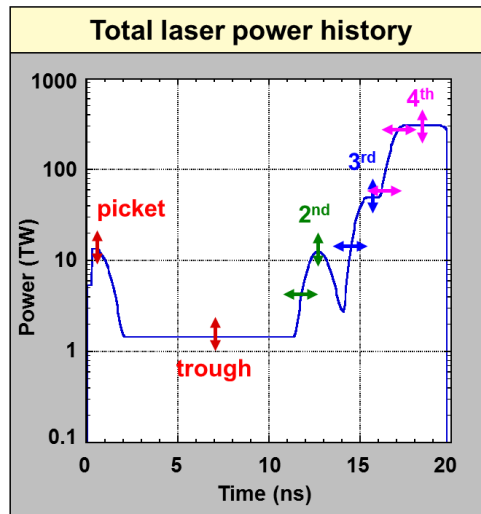


Indirect-drive (NIF)

Laser produces
x-rays inside a
hohlraum, which
irradiate the
capsule



ICF implosions result in two phases of nuclear burn: shock burn and compression burn



Hydro assumptions can break down during the shock-convergence phase

Mainline ICF simulations are made with average-ion hydrodynamic approximation

Single-fluid model

$$\frac{\partial \rho}{\partial t} + \nabla(\rho \mathbf{v}) = 0$$

$$\rho \left(\frac{\partial \mathbf{v}}{\partial t} + \mathbf{v} \cdot \nabla \right) \mathbf{v} = \nabla \mathbf{J} \times \mathbf{B} - \nabla P + \frac{\rho}{m} \mathbf{F}$$

$$\frac{m}{ne^2} \frac{\partial \mathbf{J}}{\partial t} = \mathbf{E} + \mathbf{v} \times \mathbf{B} - \frac{1}{en} \mathbf{J} \times \mathbf{B} + \frac{1}{en} \nabla P_e - \eta \mathbf{J}$$

Averaged quantities over all species

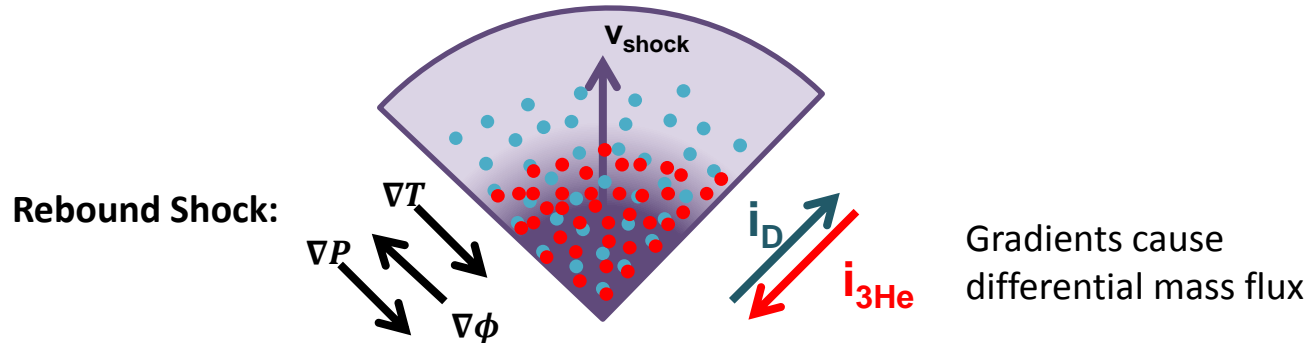
$$\rho = \sum n_i m_i + n_e m_e$$

$$P = P_i + P_e$$

$$\mathbf{v} = \frac{1}{\rho} (n_i m_i \mathbf{v}_i + n_e m_e \mathbf{v}_e)$$

$$\mathbf{J} = en(\mathbf{v}_i - \mathbf{v}_e)$$

Strong gradients in pressure, electric potential or temperature can cause species separation



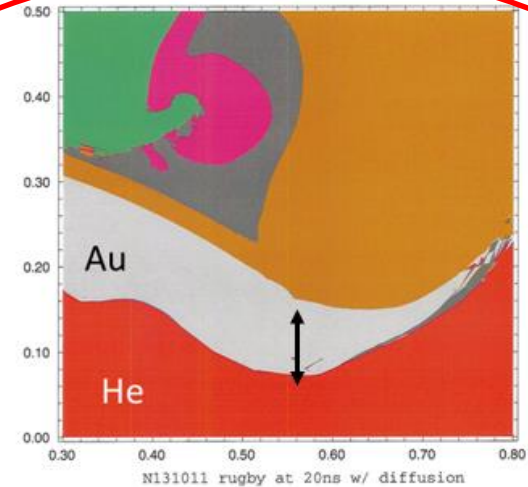
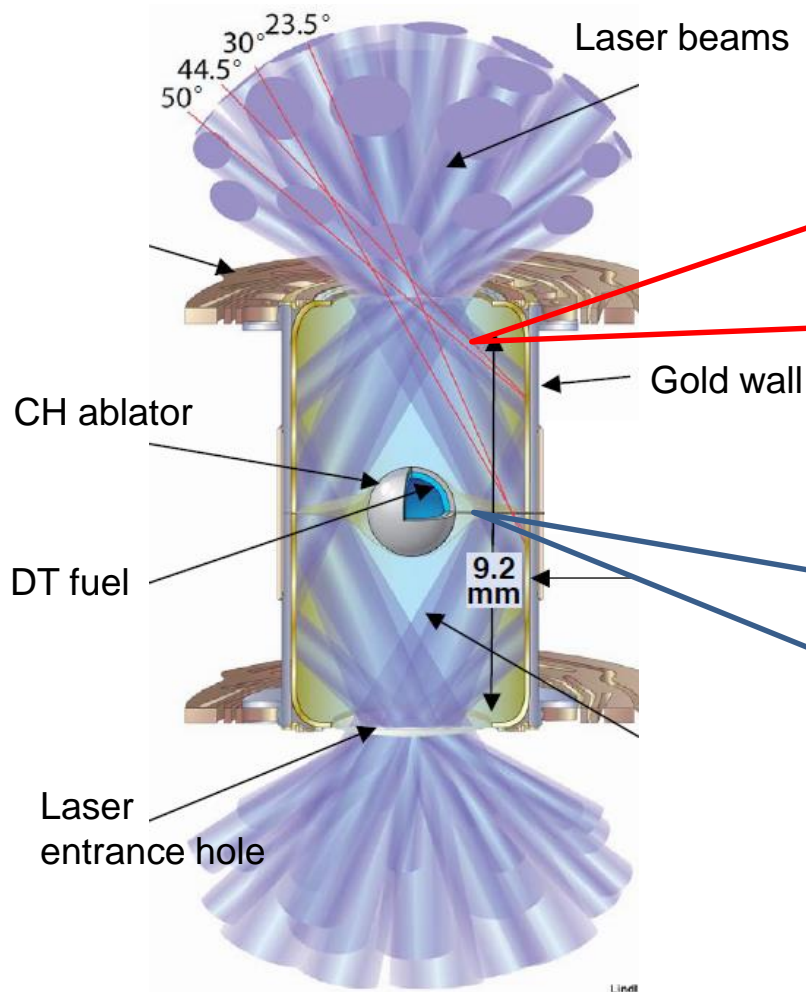
$$\text{Mass flux } i_D = -\rho D \left(\underbrace{\frac{d\alpha}{dx}}_{\text{Classical diffusion}} + \underbrace{k_p \frac{d \ln P}{dx}}_{\text{Baro-diffusion}} + \underbrace{k_E \frac{e \nabla \Phi}{k_B T}}_{\text{Electro-diffusion}} + \underbrace{k_T \nabla \ln T}_{\text{Thermo-diffusion}} \right) = -i_{3\text{He}},$$

$$\alpha = \rho_D / \rho_{tot} \sim f_D$$

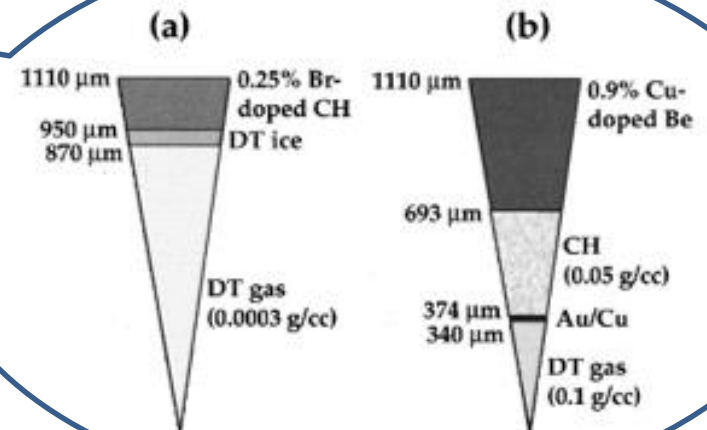
These effects will impact DT as well as D³He.

Interfaces are present in hohlraum and capsule of a typical indirect-drive ICF target

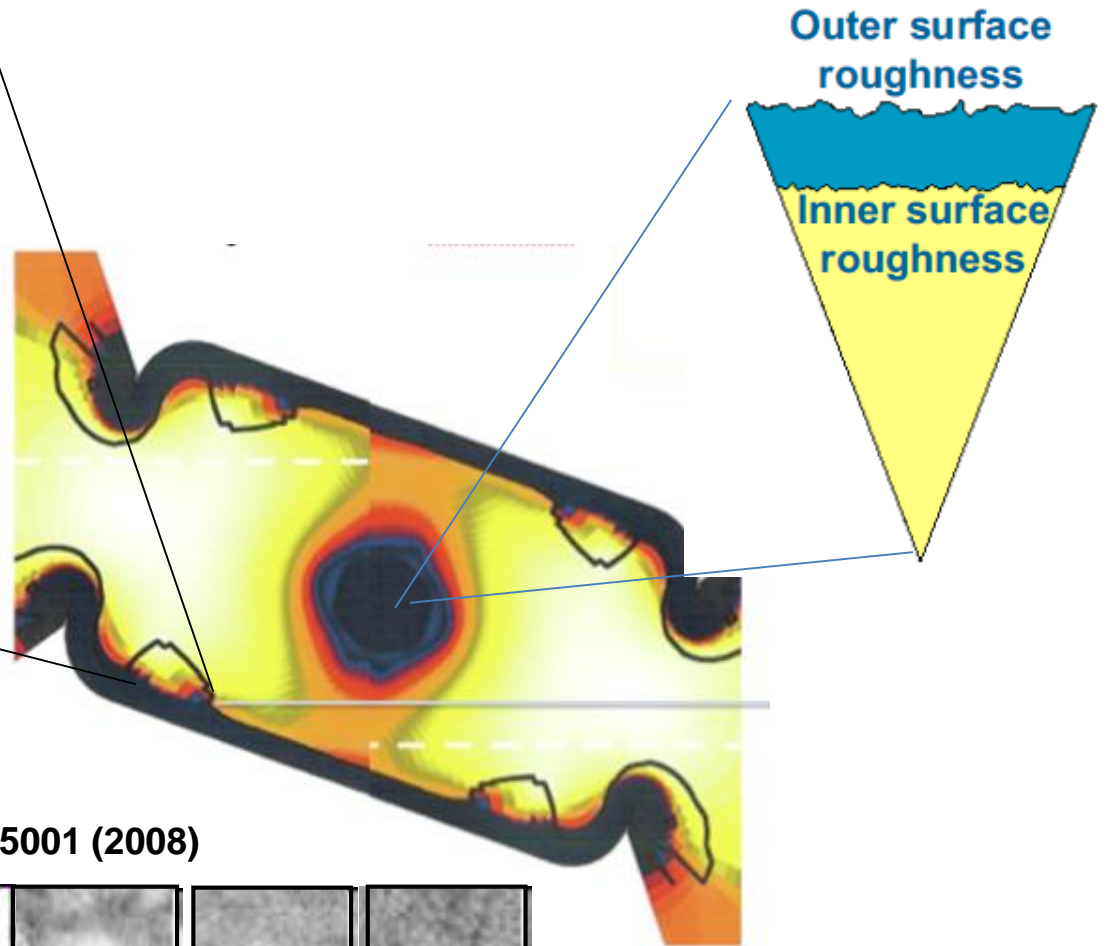
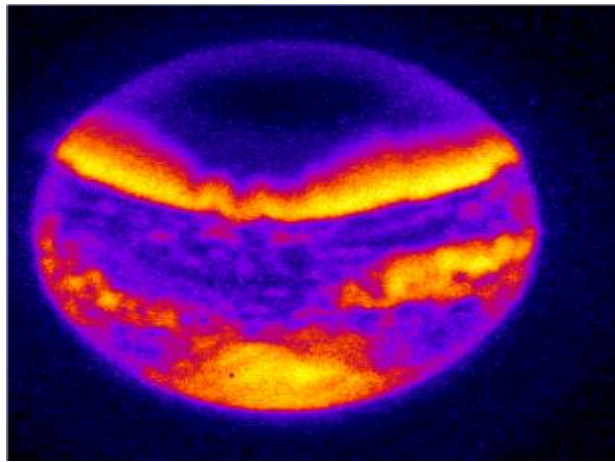
NIF Hohlräum



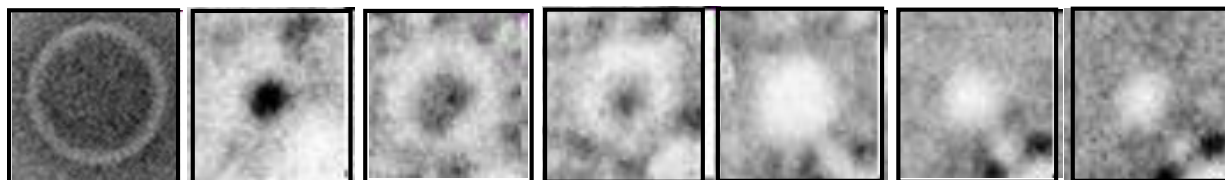
(courtesy of P. Amendt)



Interfaces are hydrodynamically unstable during laser-plasma interactions and ICF implosions, and are affected by fields



C. K. Li *et al*, Phys. Rev. Letts. 100, 225001 (2008)



0.0 ns 0.8 ns 1.2 ns 1.4 ns 1.6 ns 1.9 ns 2.1ns

Time-gated, monoenergetic proton radiography offers unique measurements of self-generated E+B fields, providing insight into ion kinetic dynamics

Imploded capsules:

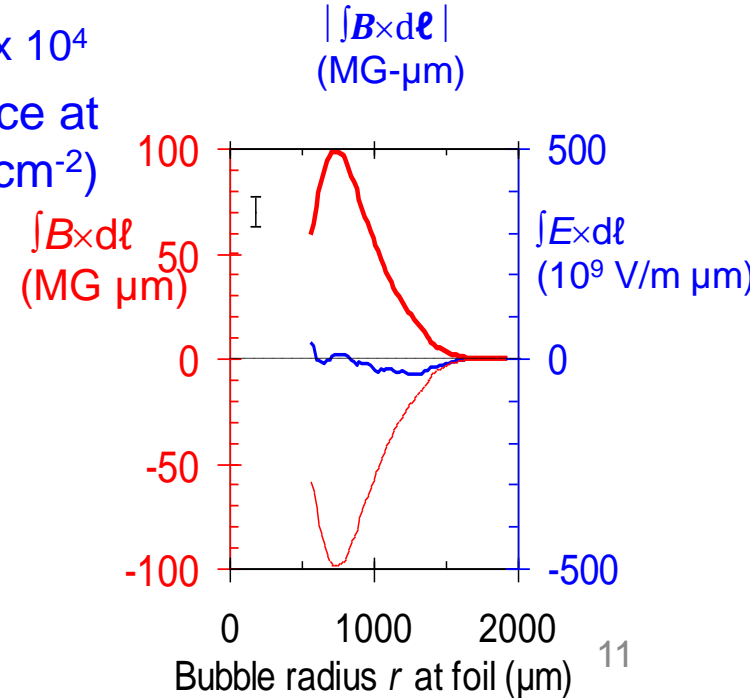
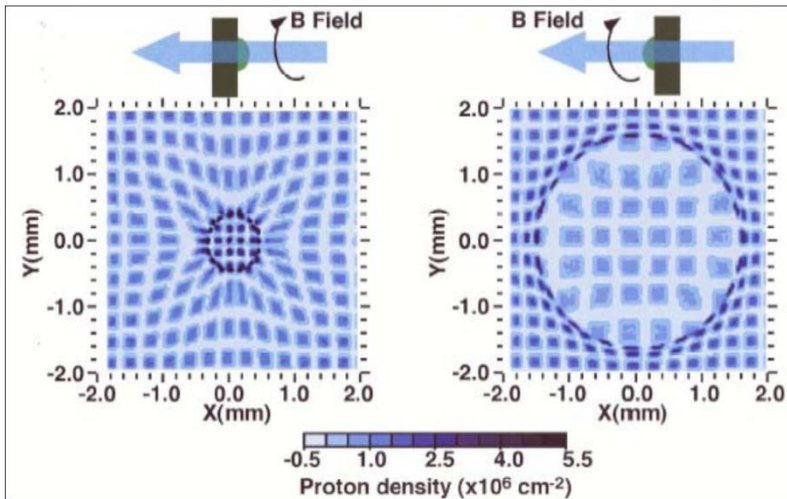
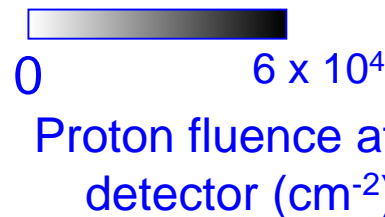
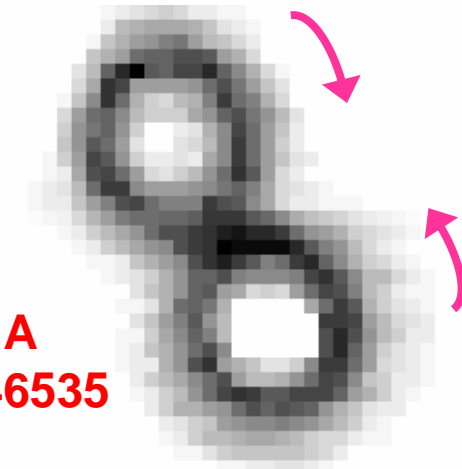
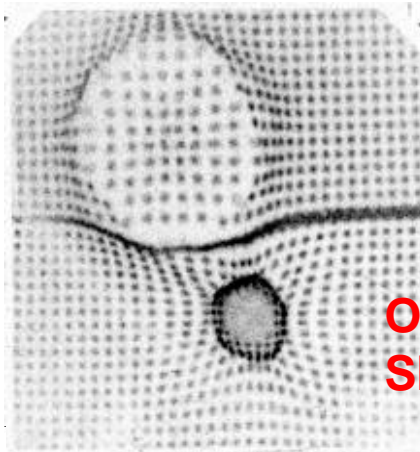
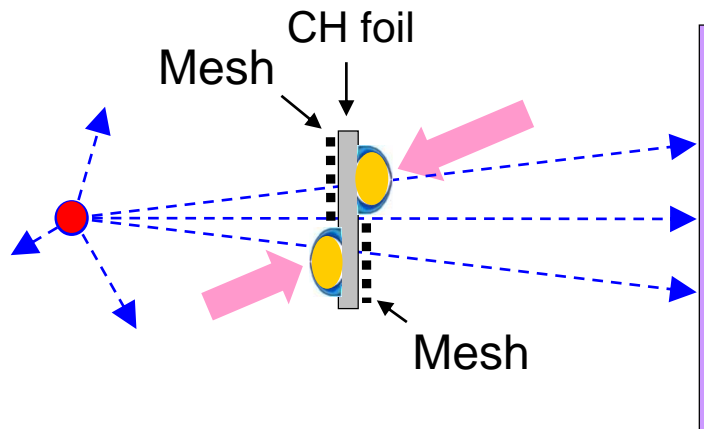
- Expansion of corona plasma \rightarrow current filaments \rightarrow B fields
- Converging shock front \rightarrow radial E fields
-

Laser-driven hohlraums

- Charge sheath formed by Ponderomotive force \rightarrow E fields
- ∇P at hohlraum wall \rightarrow E fields
- Diffusive mix at interfaces \rightarrow ambipolar E fields
- $\nabla n \times \nabla T$ around laser spots \rightarrow B fields
- Hydro unstable interfaces \rightarrow RT induced B fields
-

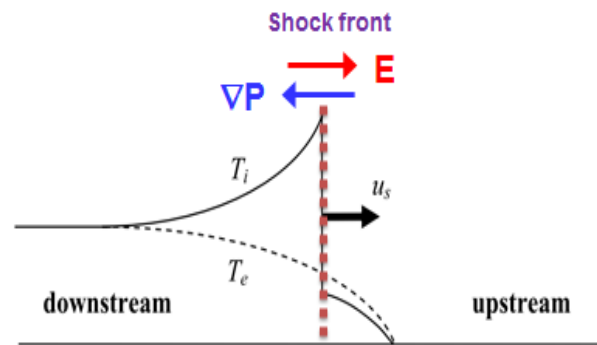
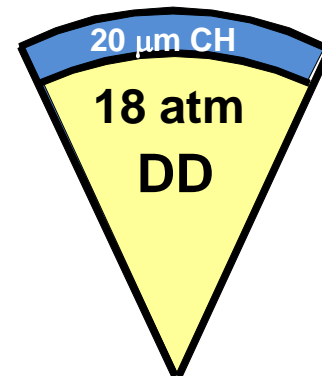
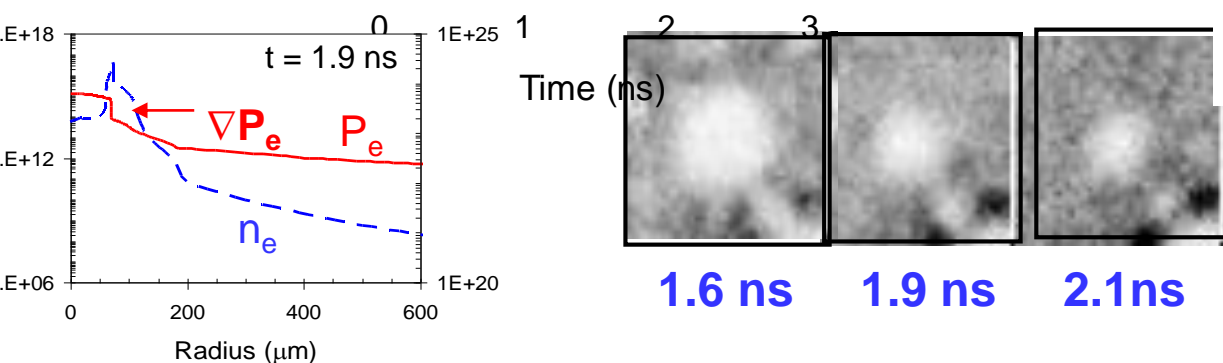
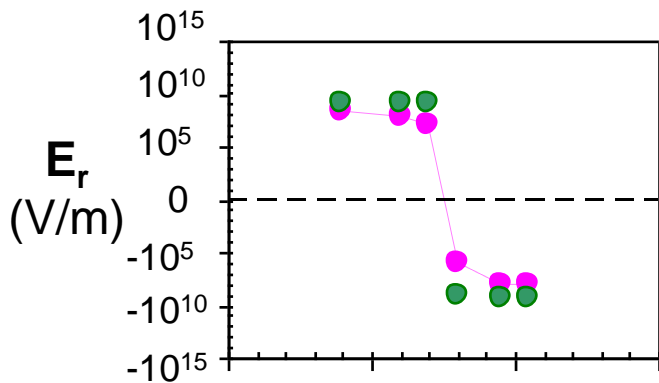
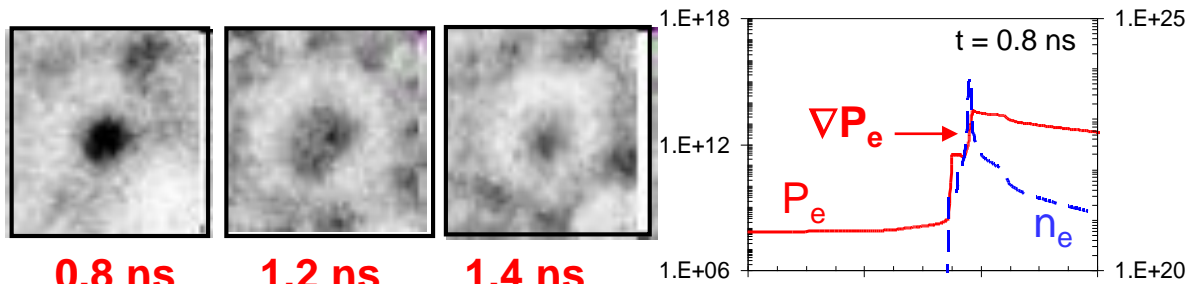
Coupled with ion-kinetic processes, self-generated fields have important effects on aspects of ICF kinetic dynamics

E & B were separated and mapped by imaging two plasma bubbles, identical except for the sign of B



Face-on radiography \rightarrow B fields
Side-on radiography \rightarrow E fields

Proton fluence focusing and its reversal indicate the direction change of self-generated radial E fields associated with moving shock front

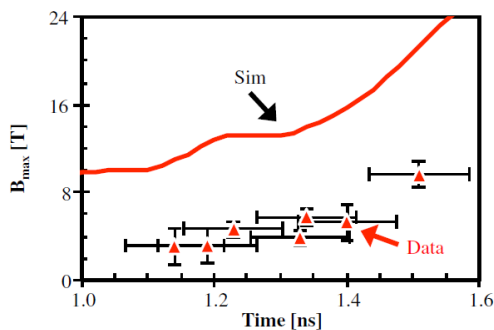
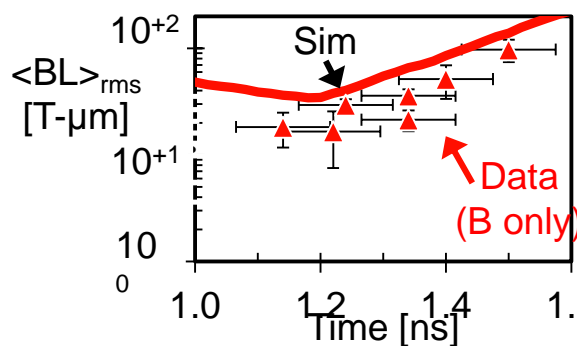
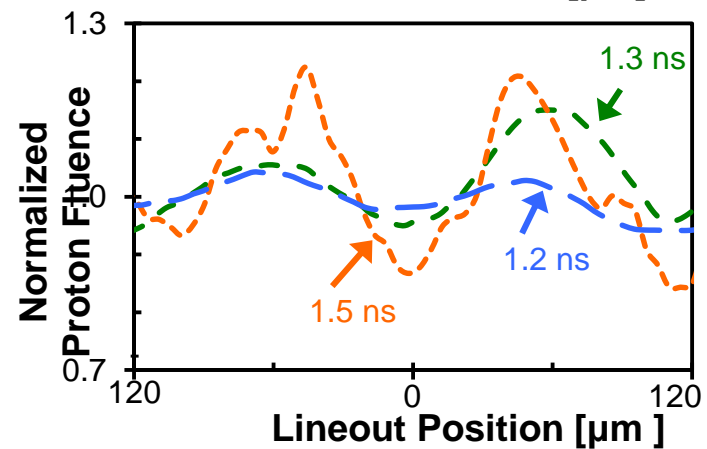
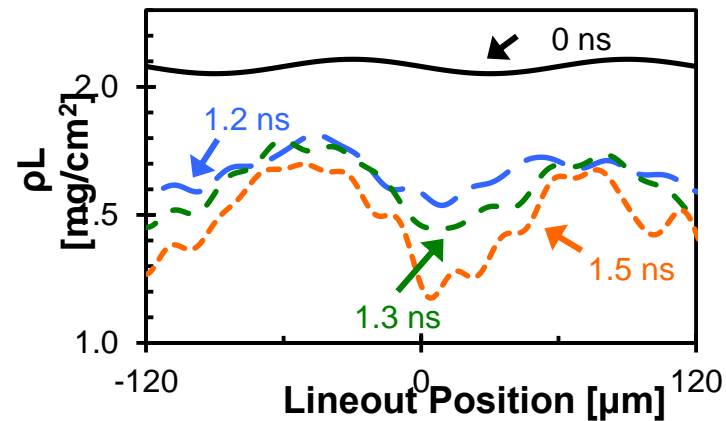
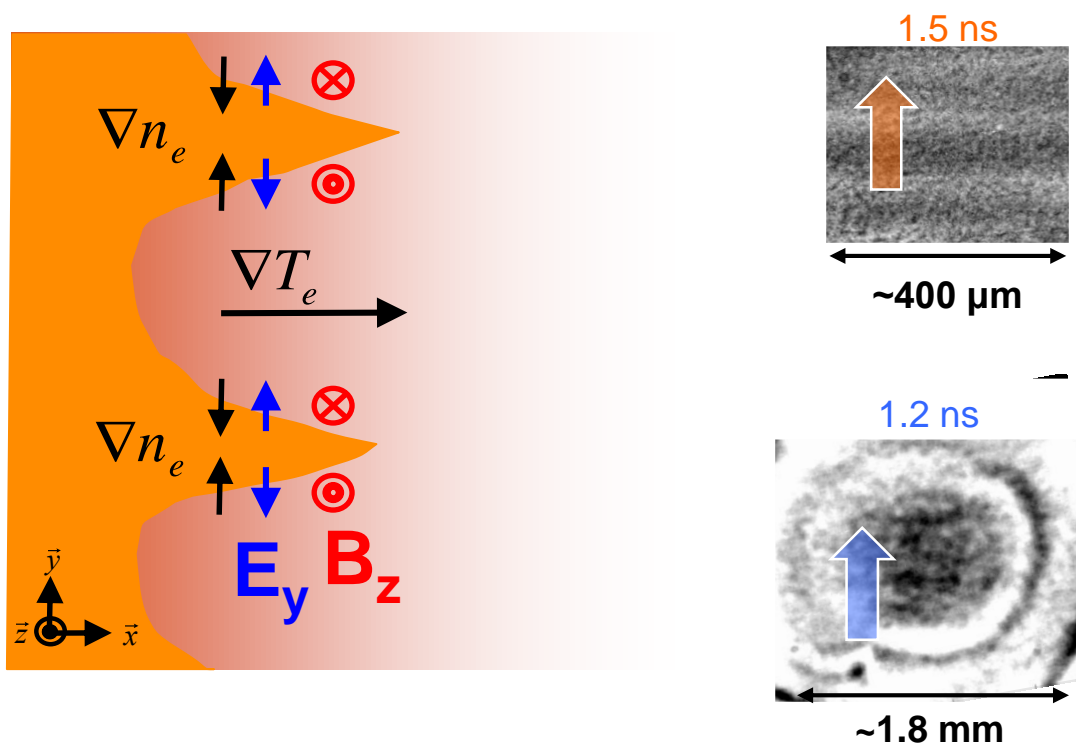


- Modify the structure of shock front
- Enhance diffusion and species separation
-

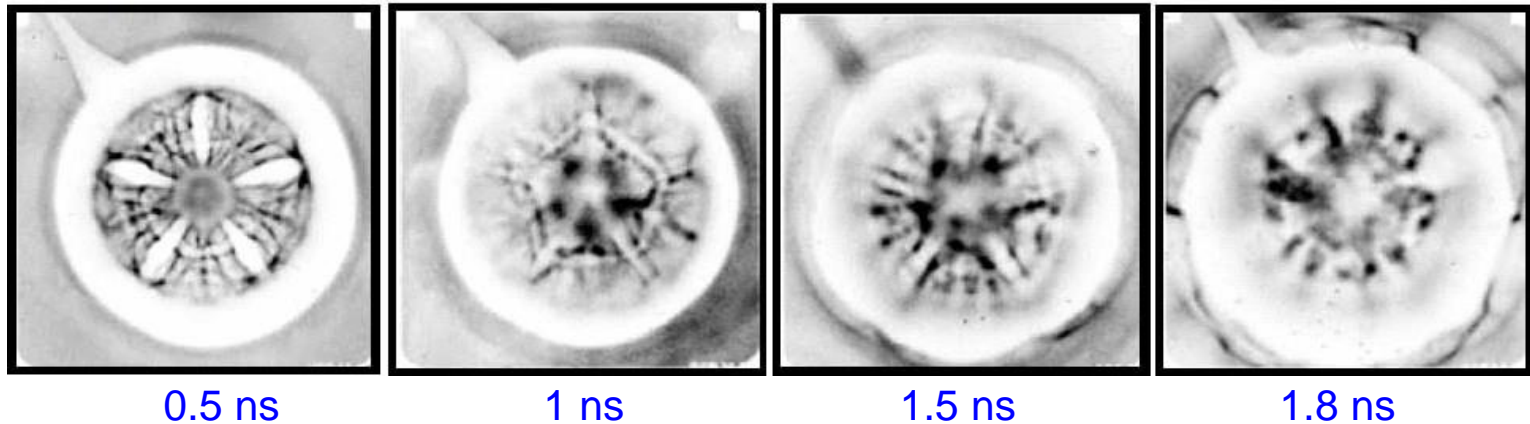
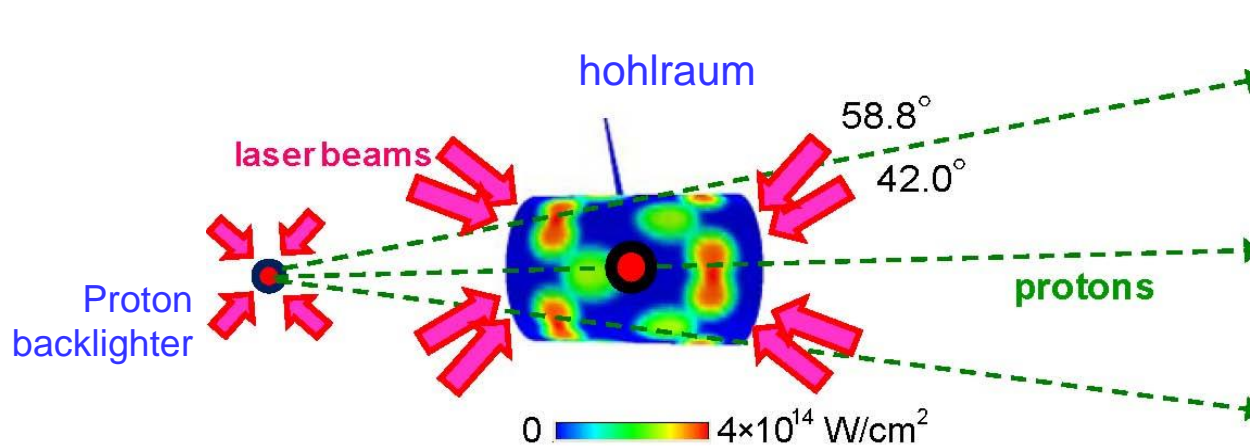
C. K. Li *et al* PRL (2008)

P. A. Amendt *et al* PRL (2010)

Proton radiography indicates the generation of B fields by Rayleigh-Taylor instabilities



Time resolved proton radiographs of indirect-drive ICF implosions at OMEGA illustrate fields and flows in hohlraums



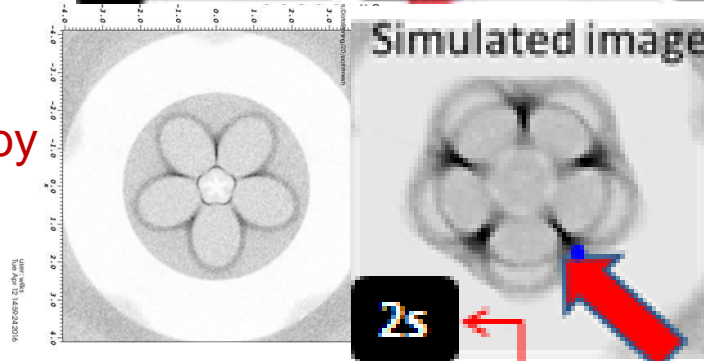
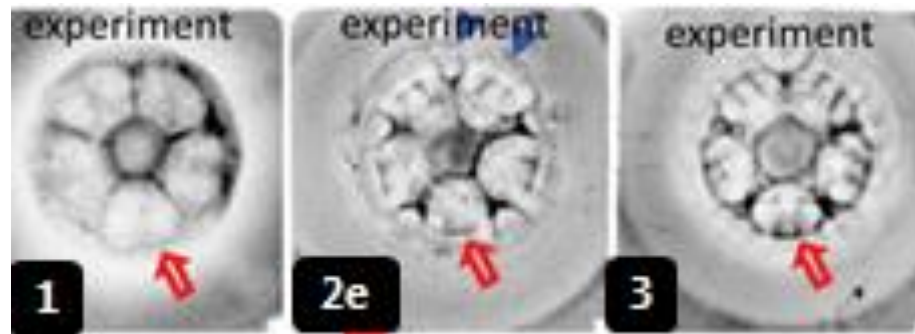
Fields are generally not included in hydro simulations

Recent simulations reproduce OMEGA experiments, indicating kinetic effects at Au/gas interfaces

OMEGA experiments

0.4 atm Neopentane

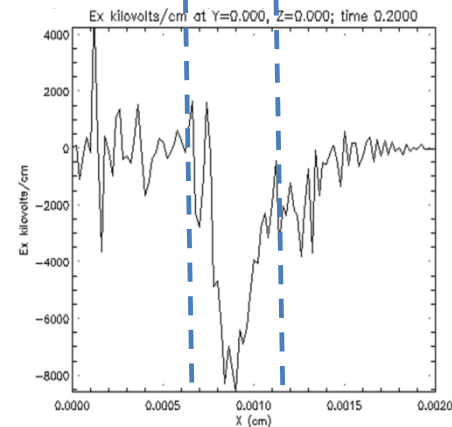
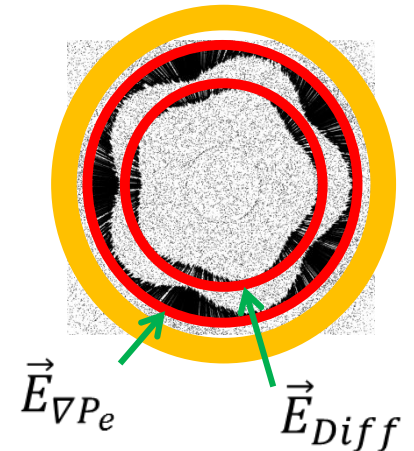
Simulation by S. Wilks



- 1** At early times, diffusion-related field and wall ∇P_e fields are in same spot.
- 2e** As time proceeds, diffusion field moves inward faster than overdense wall.
- 2s** Simulations with both E-fields included can capture these features.
- 3** Once laser is off, diffusive field dominates.

- Modify the structure of interface
- Enhance diffusion, species separation

On axis view of \vec{E}



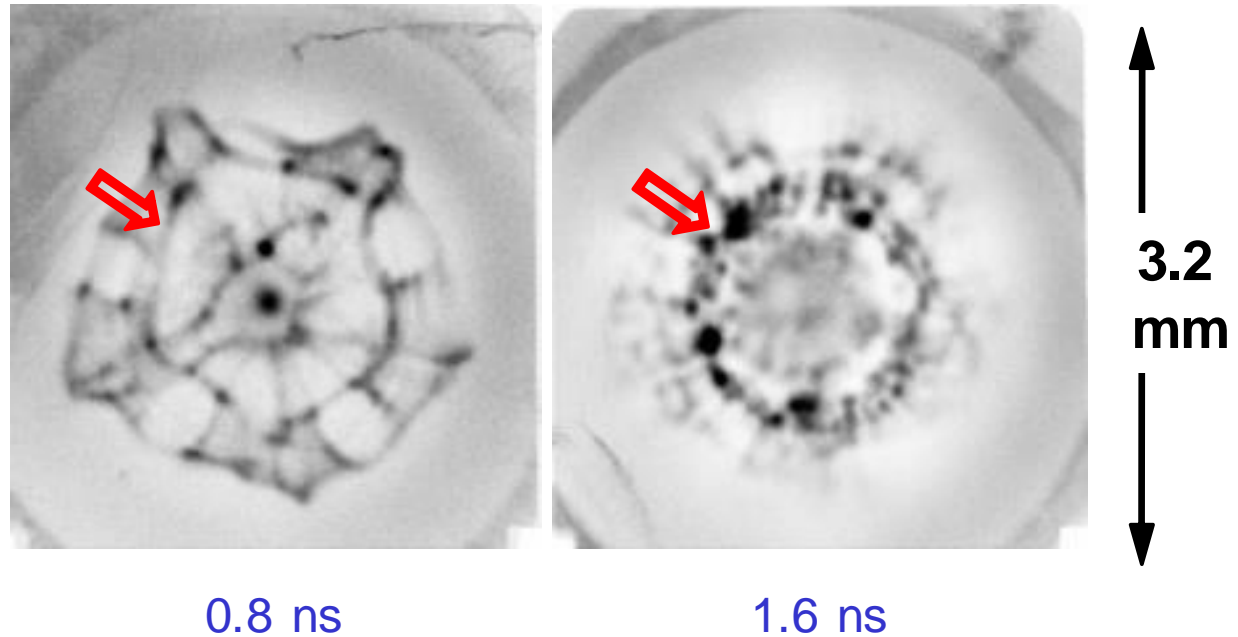
Interpenetration occurs due to the classical Rayleigh-Taylor instability as the lighter, decelerating ionized fill gas pushes against the heavier, expanding gold wall blow-off

Rayleigh-Taylor growth

$$\gamma_{\text{RT}} = \sqrt{2\pi A_t a k}$$

@ Au-Gas interface

$$A_t = \frac{\rho_2 - \rho_1}{\rho_2 + \rho_1} \approx 0.54$$



The consequence is a reduced benefit of the gas fill because the enhanced interpenetration (or mixing) between the Au blow-off and the gas plasma leads to high-Z material stagnating earlier in the hohlraum interior

Summary

Peak values are $E \sim 10^9 \text{ V m}^{-1}$ and $B \sim 10^6 \text{ gauss}$ in different ICF implosion scenario

- **ICF capsule implosions**
- **Laser-irradiated hohlraums**
- **Quantitative measurements of the effects of such fields on ICF implosions are difficult ongoing undertaking**