

(Magnetized) collisionless ion- gas/plasma interactions

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modelling the streaming instability in the laboratory?

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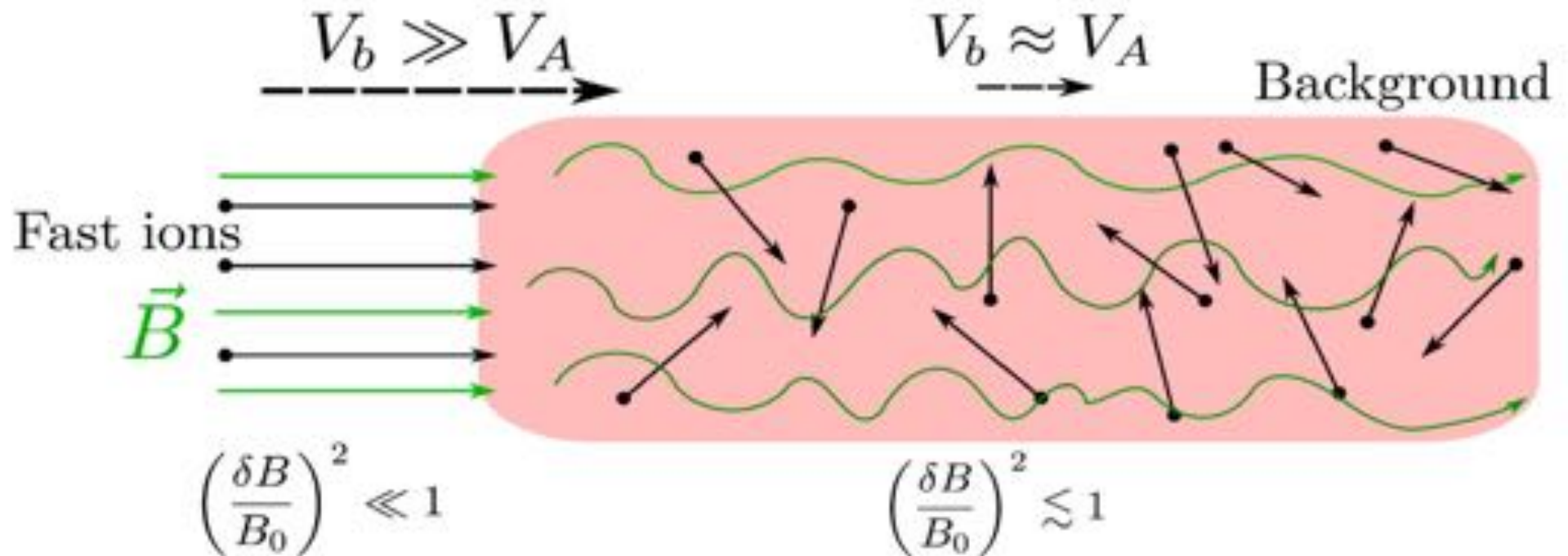


Schematic of the streaming instability

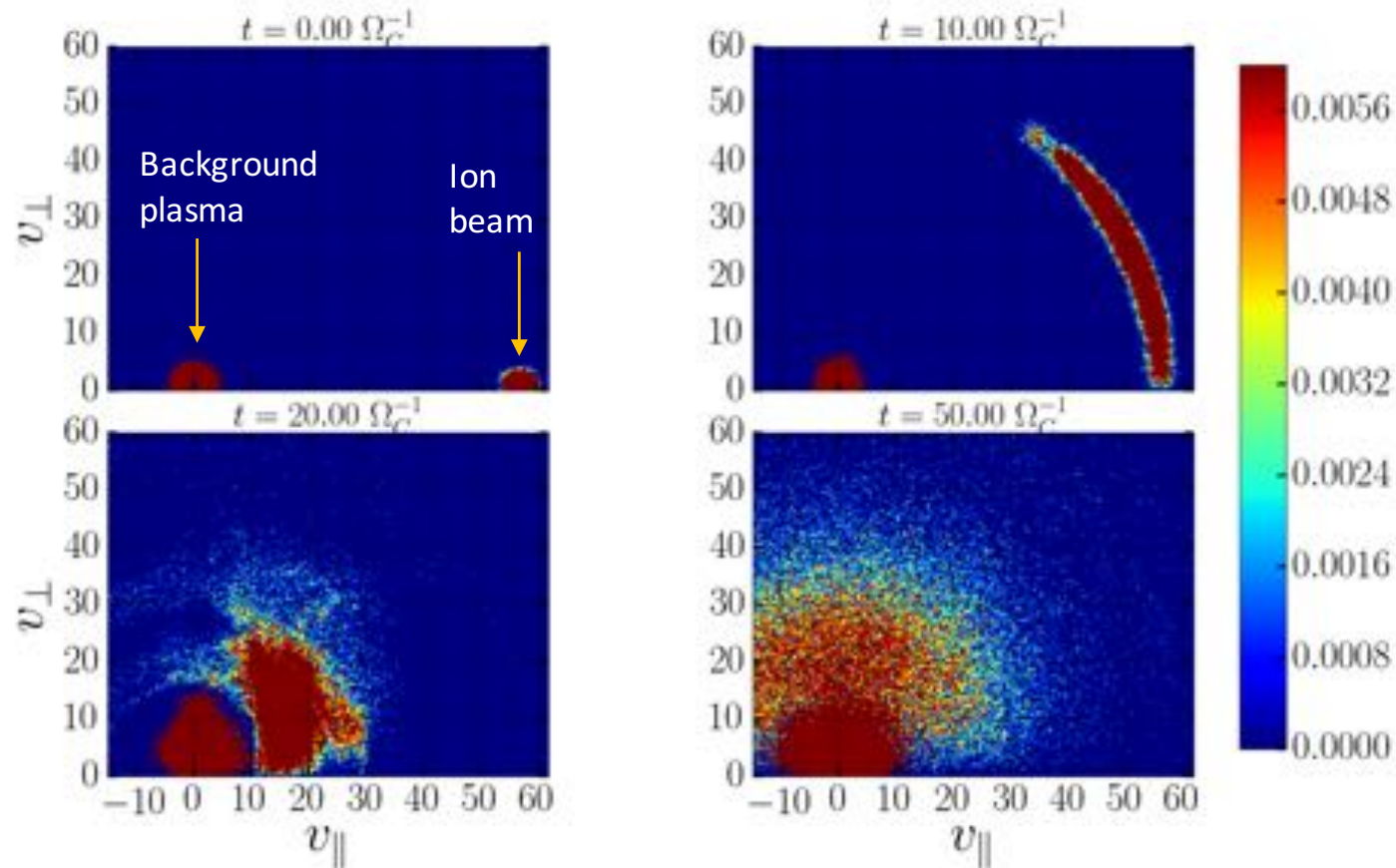
- EM ion instability ubiquitous in space plasmas
- Generates MHD waves & turbulence
- Suspected to play a key role in “anomalous” ionization of ISM, particle acceleration at shocks, ...

Growth rate $\sim \Omega_{ci}$

Length scales $\sim c/\omega_{pi}$



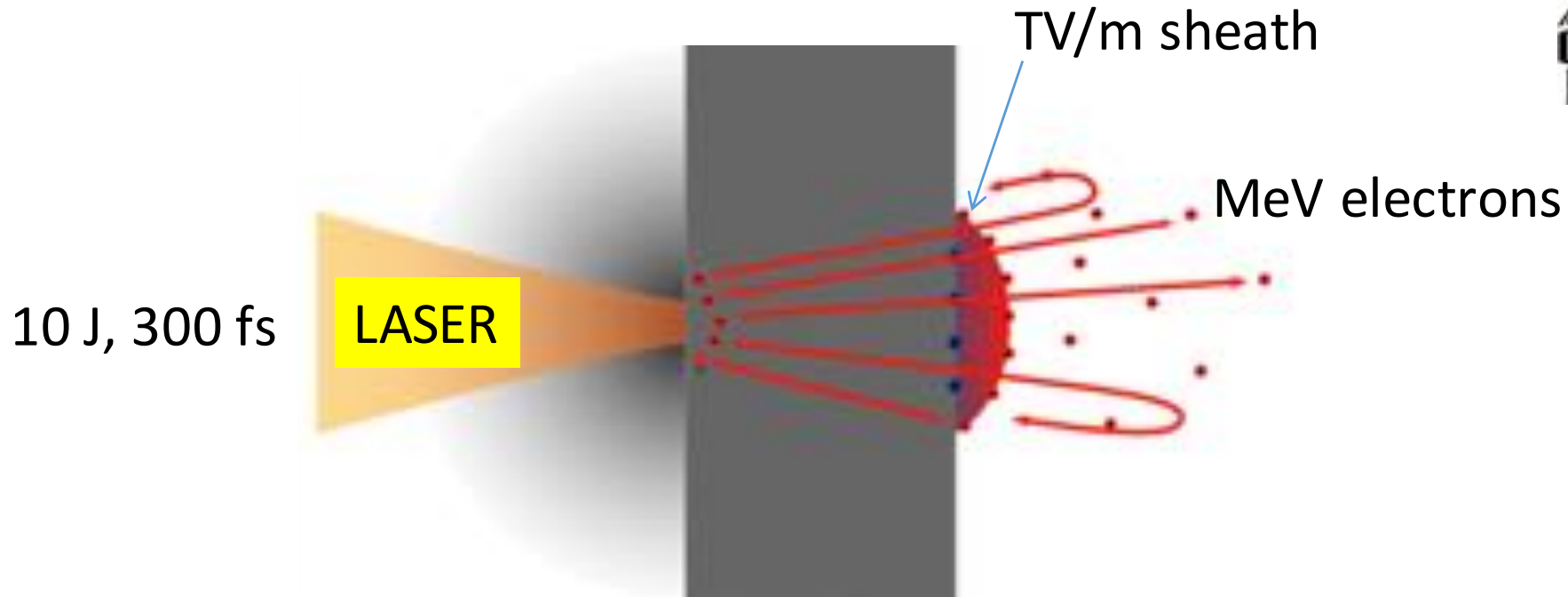
Hybrid (Heckle code) illustrative simulation



- initial streaming proton at $v_{||}=55$ (in units of Alfvén speed, i.e. 300 keV)
- 20 T magnetized plasma

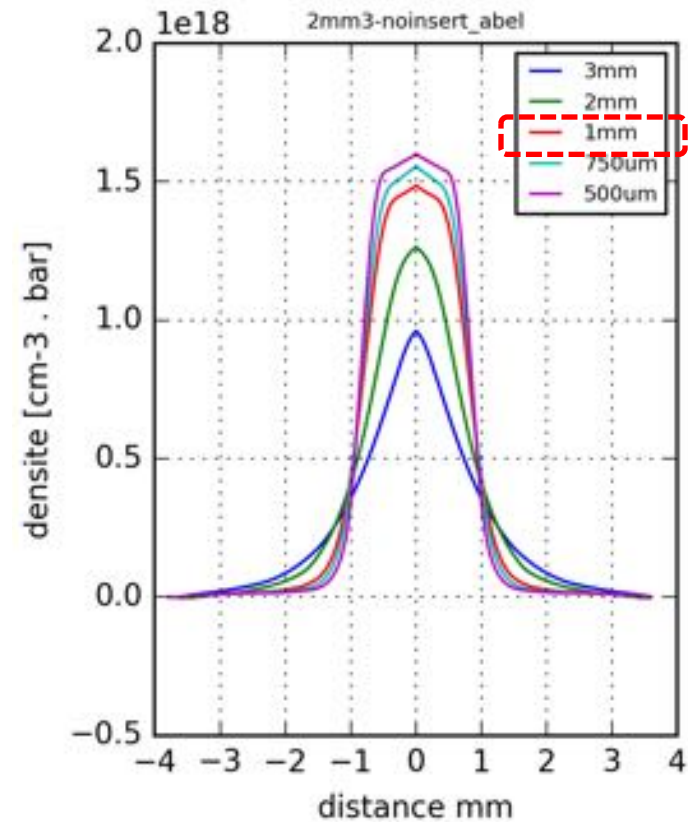
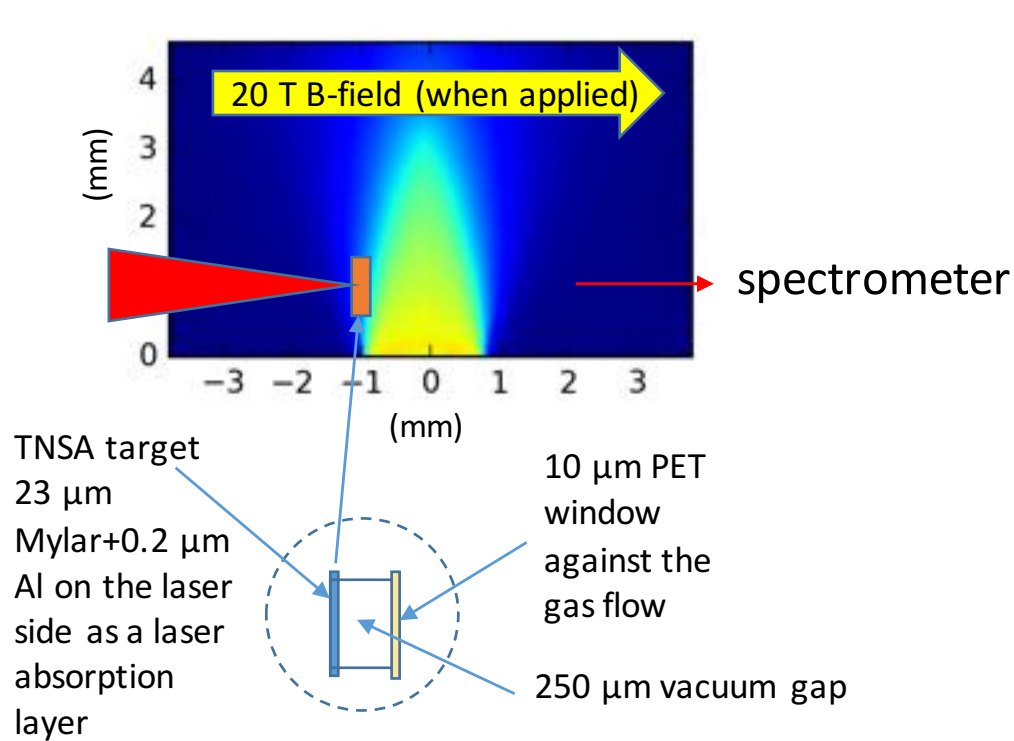
→ progressive thermalization of the directed ion beam

We exploit ultra-intense lasers electrostatic acceleration on solids (“TNSA”) to generate the fast ions

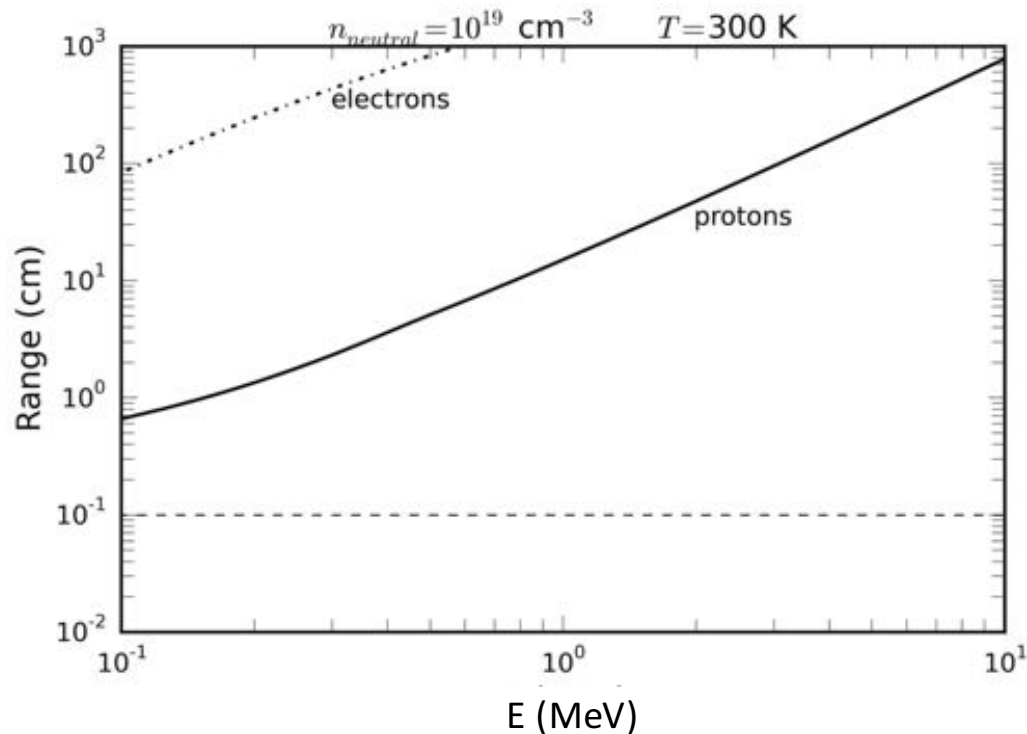


- Maximum proton energy is near 10 MeV.
- Divergent beam.

For the plasma medium, we use a gas jet



For these densities ($\sim 10^{19} \text{ cm}^{-3}$), we are in a collisionless configuration



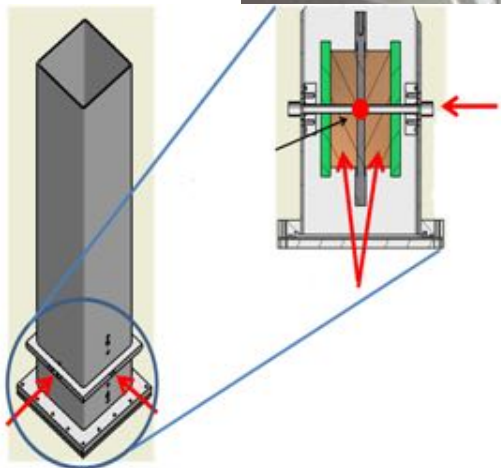
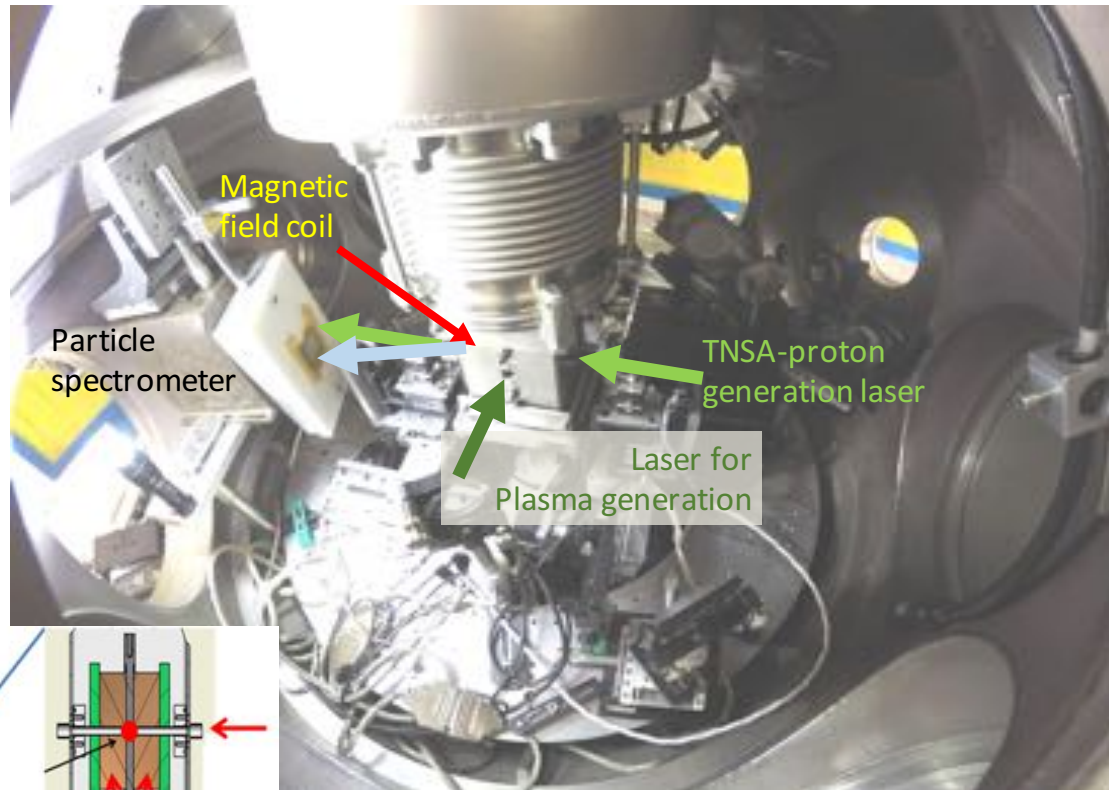
For a 1 MeV proton:

Gyro-period $\sim 0.5 \text{ ns}$

Mean free path $>$ system size (few mm)

$v \sim 3 \times v_{\text{Alfven}}$

Our method for large-scale, uniform external magnetization

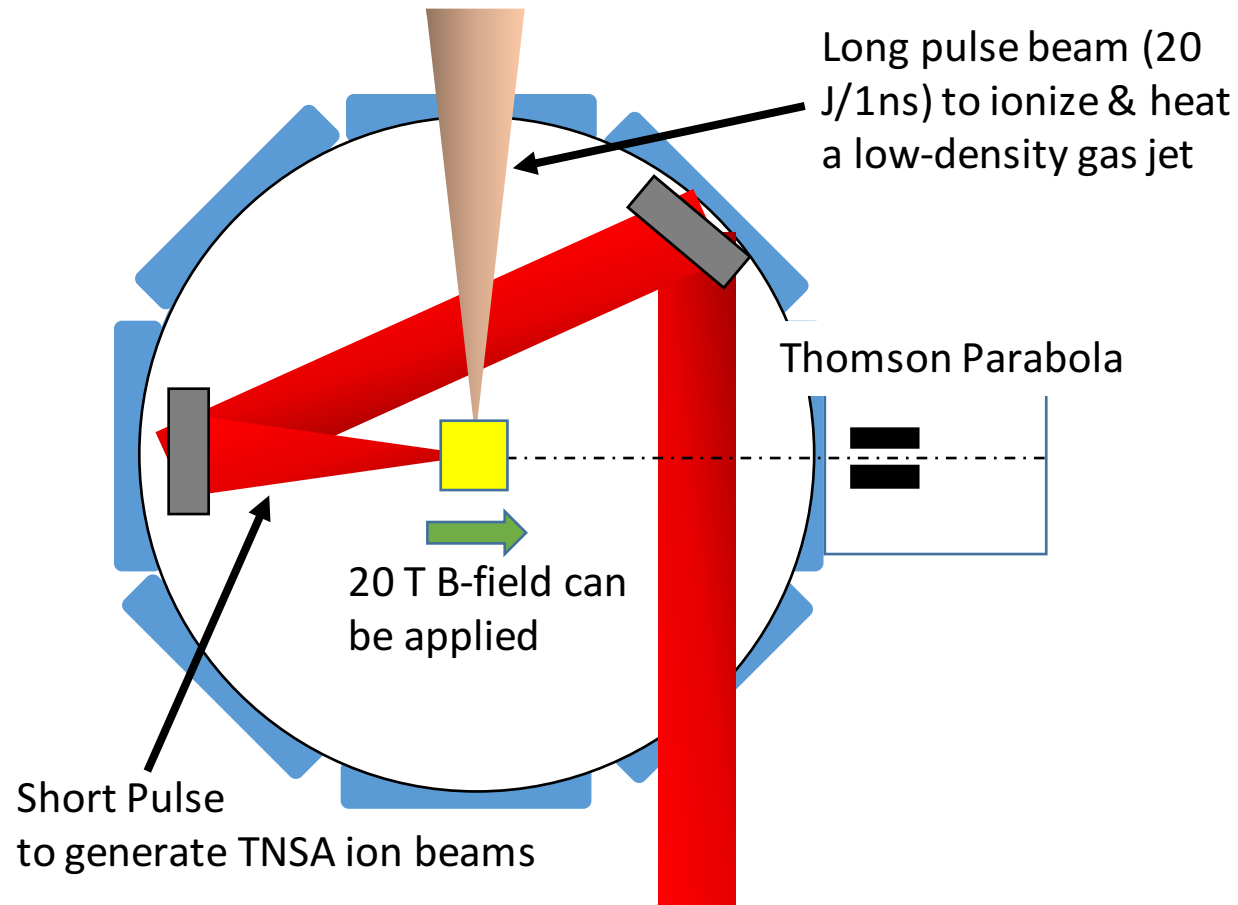


**Typical
plasma scales**
mm, ps

B-field scales
cm, μ s

[B. Albertazzi et al., Rev. Sci. Instru 84, 043505 (2013)]

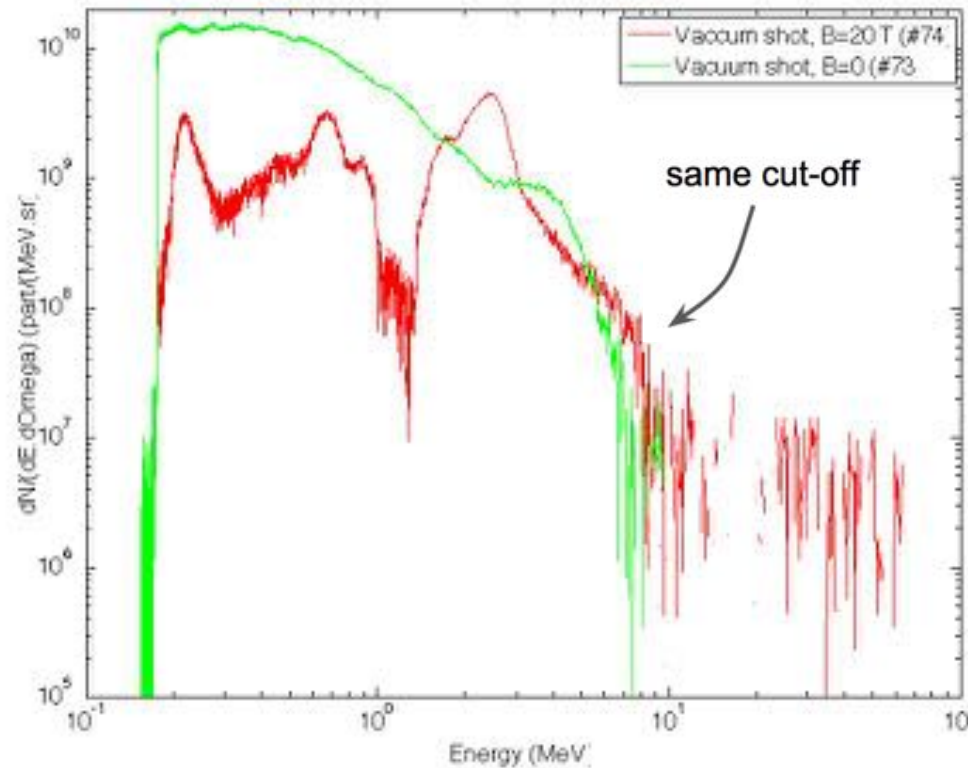
Overall set-up



First vacuum test shots : Immersing the setup in 20 T B-field



→ does not affect the proton maximum energy from the TNSA target, but modulates the spectrum



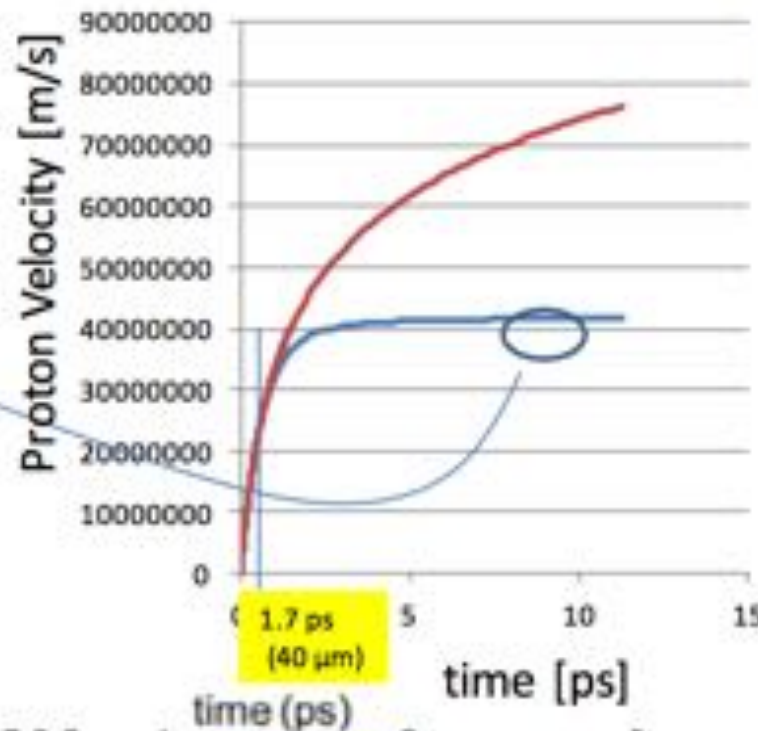
Magnetization would impact the protons if the B-field magnetizes the electrons before acceleration is saturated



Two phases in TNSA:

- 1) Isothermal
- 2) Adiabatic cooling of the electrons

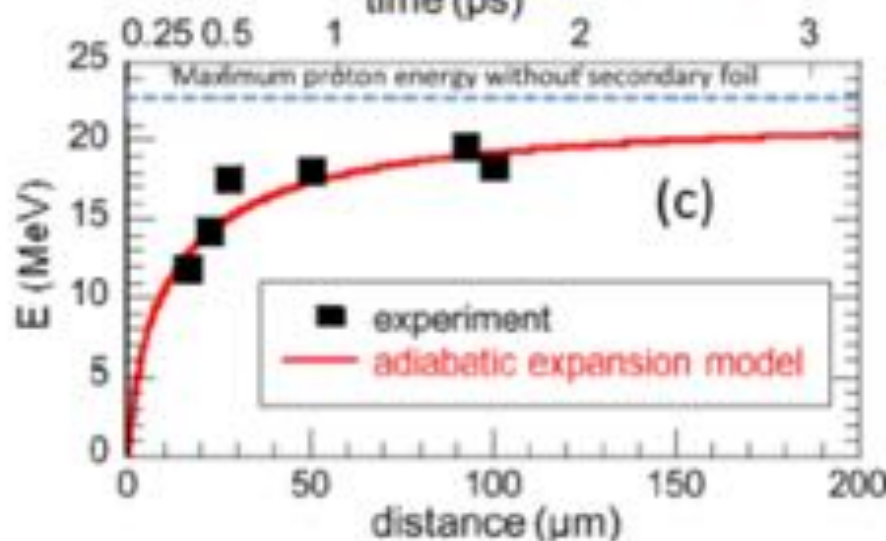
$$v_{final} \cong 2c_s \ln(0.32 x_L / \lambda_D + 4.2)$$



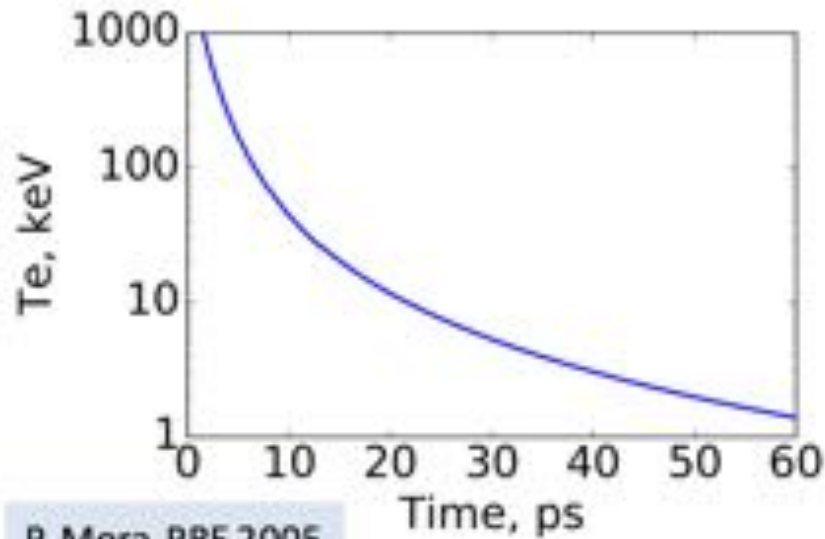
P. Mora, PRE 2005

- adiabatic
- isothermal

Corroborated by our measurement of proton energy saturation using similar parameters [S.N. Chen, PoP 2014]



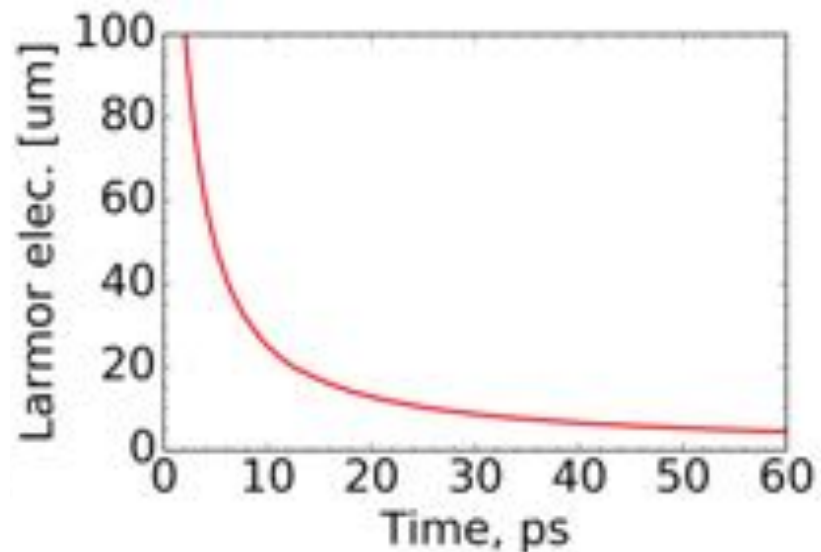
Before 2 ps, the electrons are indeed not magnetized, hence no impact of the B-field on acceleration



P. Mora, PRE 2005

$$\text{For: } t < \tau_{XL} \\ \rightarrow T_e \sim T_{e0}$$

$$\text{For: } t > \tau_{XL} \\ \rightarrow T_e(t) \propto T_{e0} \left(\frac{\tau_{XL}}{t} \right)^2$$



When the TNSA beam reaches > 100 μm , the electrons have fully cooled down.



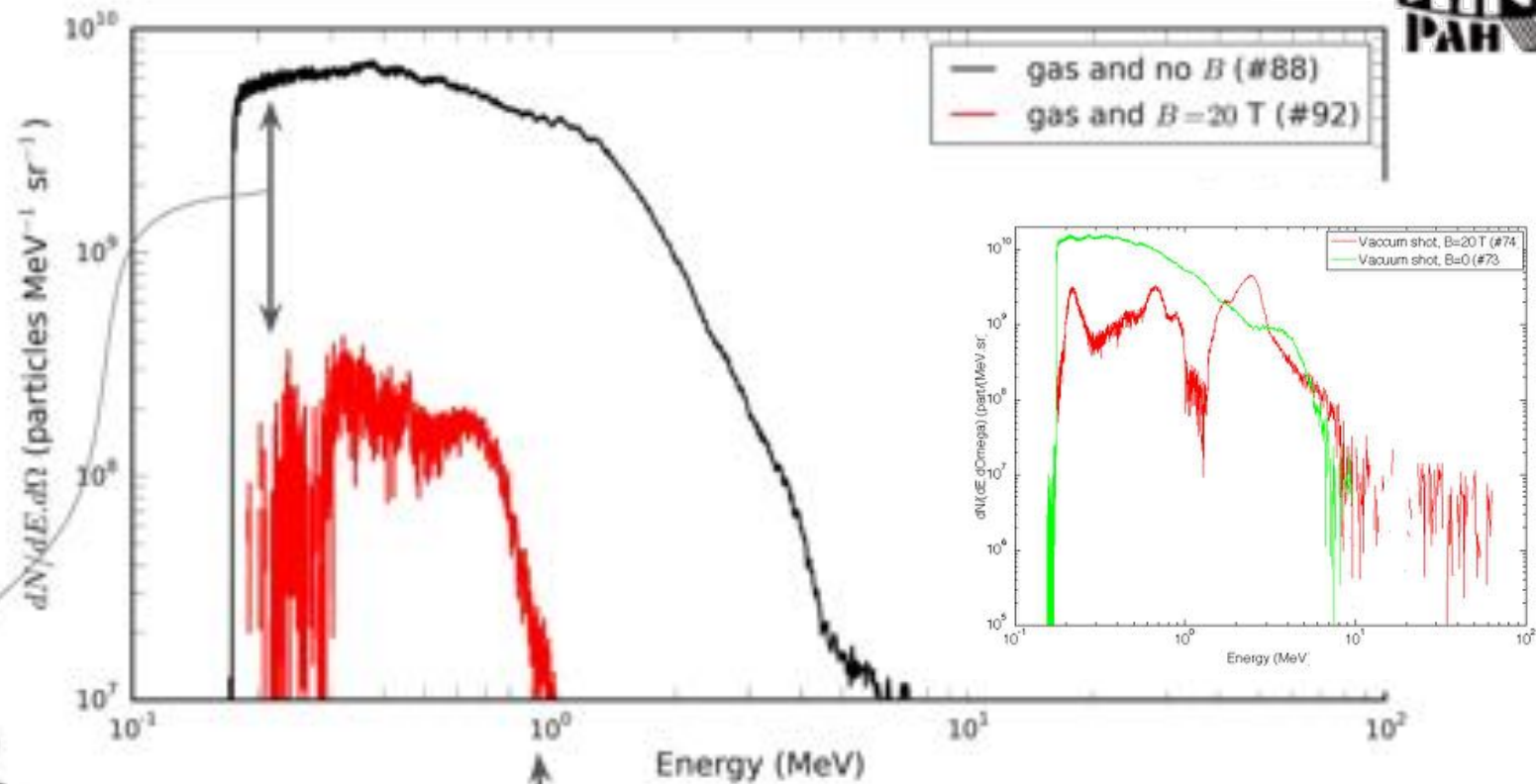
By this time their gyroradii are small compared to the system size.

When injected in the *magnetized* (non pre-ionized) gas, a huge downshift in proton energy is observed



Ne gas @ 15 bars
 → $\sim 2 \cdot 10^{20}$ at/cc

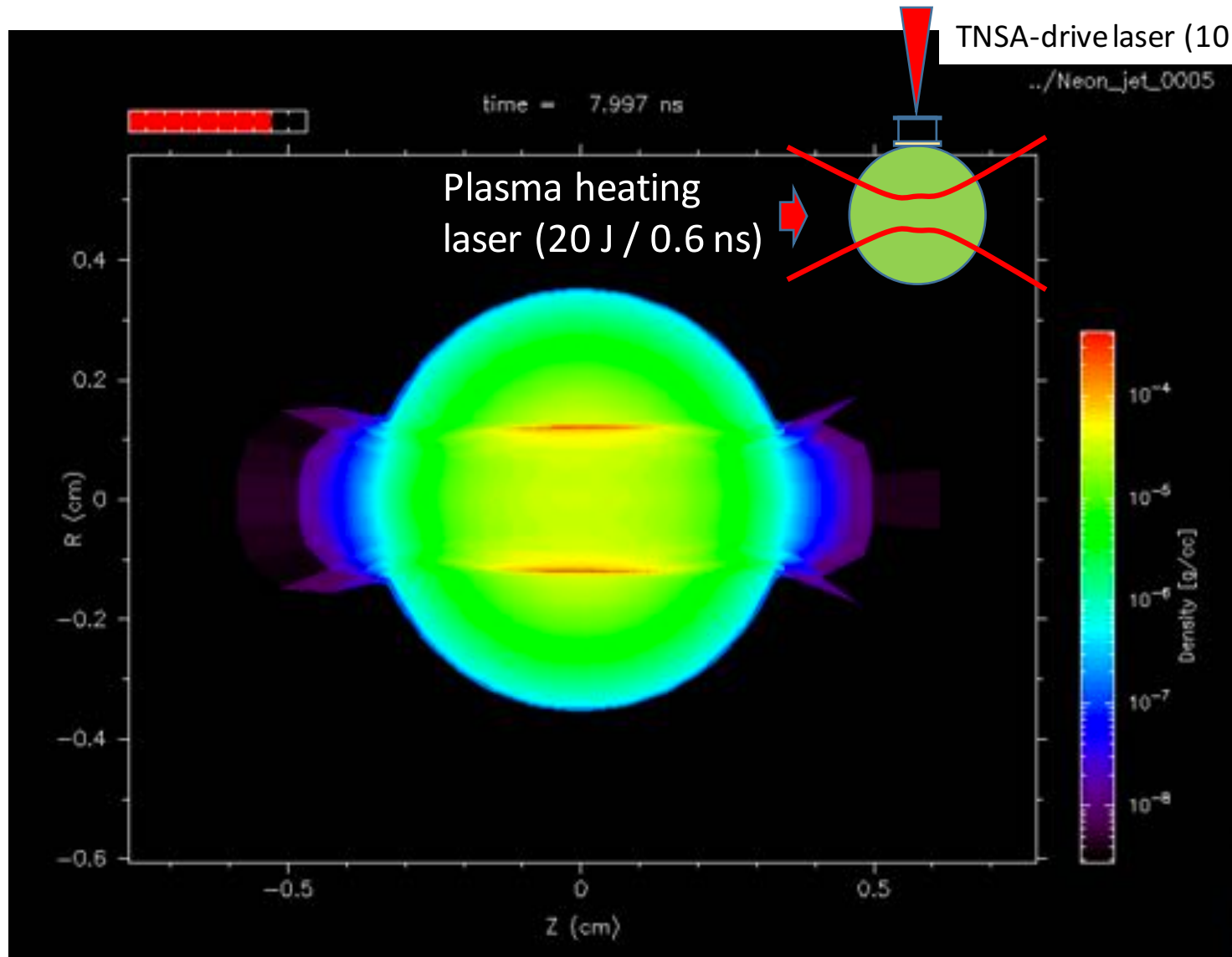
nanosecond-laser OFF



Shift cannot be due to magnetic field modulation only.
 Extra scattering?

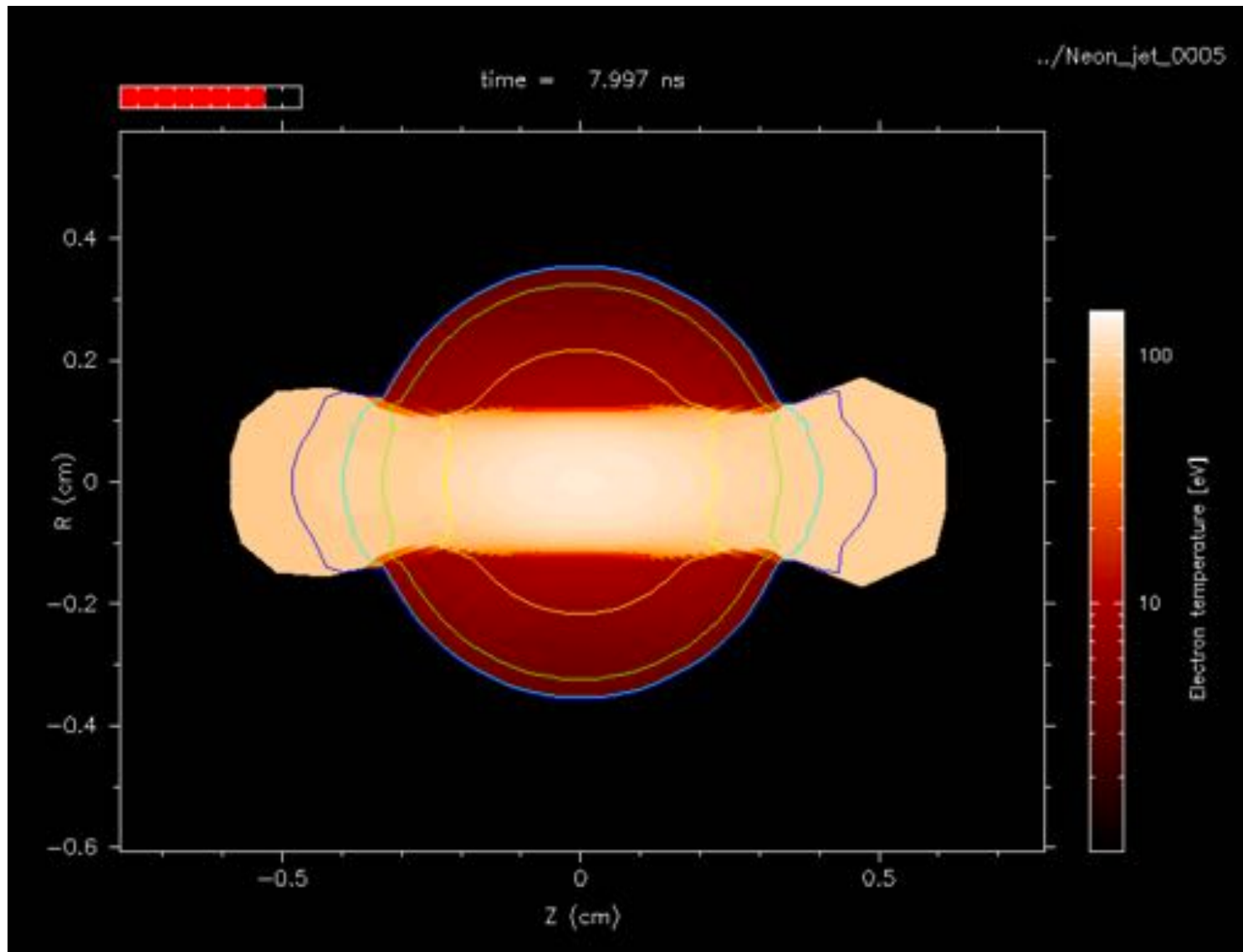
Large change in the energy cut-off

Now heating by the long-pulse, quite homogeneously



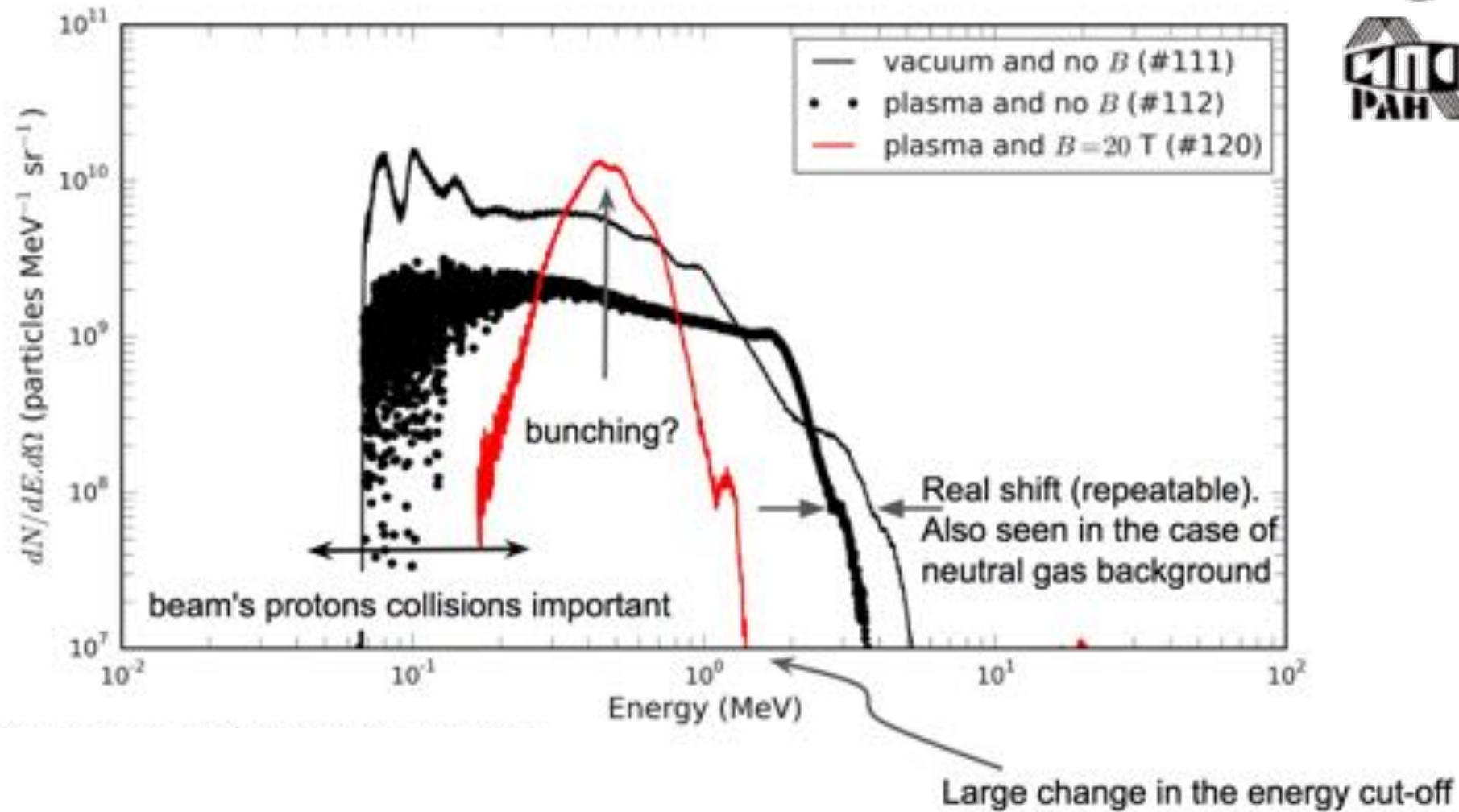
Simulations
using the
DUED code
(S. Atzeni),
Ne gas

The long-pulse induces a quite homogeneous heating



Simulations
using the
DUED code
(S. Atzeni)

Proton beam in plasma: also large downshift



Wrap-up



Beam through

- vacuum **with** magnetic field
 - modulation of the spectrum
 - no change in cut-off energy
- "neutral" gas/plasma **without** magnetic field
 - small downshift in cut-off energy
- "neutral" gas or heated plasma **with** magnetic field
 - beam-background particle collisions mostly negligible → **collective effects** (?)
 - large modification of spectrum
 - much smaller energy cut-off